

ECE 590.01 Quantum Engineering with Atoms

Spring 2020

Mid-Term Exam #1 – Revised 3/2/2020

Due on Friday 3/6/2020 12:00 pm EST

Instructor: Geert Vrijsen and Jungsang Kim

Name: _____

I hereby certify that I worked on this mid-term exam strictly abiding by the rules set forth by the instructor, as indicated below. I understand that the violation of these rules will be considered an academic misconduct, and will be subject to punishment. I also certify that I participated in the exam following the academic integrity and honesty anticipated for all Duke students.

Exam Rules:

1. I did not utilize any other documents, in paper, electronic or in any other format, other than the following allowed documents: the recommended textbook for the class (Cohen-Tannoudji, Diu and Laloe, Quantum Mechanics), classroom material provided through the website, notes I have taken in the classroom, and the homework problems.
2. I did not utilize any other resources to solve the exam problems, including any other textbooks or papers, homework/exam problems/solutions from similar classes taught in the past at Duke or elsewhere, or information available through web searches.
3. I worked on the problems by myself, and did not discuss the problems with anyone else other than the instructor; including my classmates, friends or colleagues, other professors, etc.

Signature _____ Date _____

2. Angular Momentum

- a. From the fundamental commutation relation $[\hat{x}_i, \hat{p}_j] = i\hbar\delta_{ij}$, prove the following relations where r is the radial coordinate in spherical coordinates (10 points)

$$[\hat{L}_i, \hat{x}_j] = i\hbar\varepsilon_{ijk}\hat{x}_k, [\hat{L}_i, \hat{L}_j] = i\hbar\varepsilon_{ijk}\hat{L}_k, \text{ and } [\hat{L}_i, f(r)] = 0.$$

- b. For the eigenstate of the \hat{L}^2 and the \hat{L}_z operator $|l = 1, m = 1\rangle$, find the probability that the measurement of \hat{L}_x would yield the values \hbar , 0 and $-\hbar$. (10 points).

Spherical Harmonics

$$l = 0$$

$$Y_0^0(\theta, \varphi) = \frac{1}{2} \sqrt{\frac{1}{\pi}}$$

$$l = 1$$

$$Y_1^{-1}(\theta, \varphi) = \frac{1}{2} \sqrt{\frac{3}{2\pi}} \cdot e^{-i\varphi} \cdot \sin \theta = \frac{1}{2} \sqrt{\frac{3}{2\pi}} \cdot \frac{(x - iy)}{r}$$

$$Y_1^0(\theta, \varphi) = \frac{1}{2} \sqrt{\frac{3}{\pi}} \cdot \cos \theta = \frac{1}{2} \sqrt{\frac{3}{\pi}} \cdot \frac{z}{r}$$

$$Y_1^1(\theta, \varphi) = -\frac{1}{2} \sqrt{\frac{3}{2\pi}} \cdot e^{i\varphi} \cdot \sin \theta = -\frac{1}{2} \sqrt{\frac{3}{2\pi}} \cdot \frac{(x + iy)}{r}$$

$$l = 2$$

$$Y_2^{-2}(\theta, \varphi) = \frac{1}{4} \sqrt{\frac{15}{2\pi}} \cdot e^{-2i\varphi} \cdot \sin^2 \theta = \frac{1}{4} \sqrt{\frac{15}{2\pi}} \cdot \frac{(x - iy)^2}{r^2}$$

$$Y_2^{-1}(\theta, \varphi) = \frac{1}{2} \sqrt{\frac{15}{2\pi}} \cdot e^{-i\varphi} \cdot \sin \theta \cdot \cos \theta = \frac{1}{2} \sqrt{\frac{15}{2\pi}} \cdot \frac{(x - iy)z}{r^2}$$

$$Y_2^0(\theta, \varphi) = \frac{1}{4} \sqrt{\frac{5}{\pi}} \cdot (3 \cos^2 \theta - 1) = \frac{1}{4} \sqrt{\frac{5}{\pi}} \cdot \frac{(2z^2 - x^2 - y^2)}{r^2}$$

$$Y_2^1(\theta, \varphi) = -\frac{1}{2} \sqrt{\frac{15}{2\pi}} \cdot e^{i\varphi} \cdot \sin \theta \cdot \cos \theta = -\frac{1}{2} \sqrt{\frac{15}{2\pi}} \cdot \frac{(x + iy)z}{r^2}$$

$$Y_2^2(\theta, \varphi) = \frac{1}{4} \sqrt{\frac{15}{2\pi}} \cdot e^{2i\varphi} \cdot \sin^2 \theta = \frac{1}{4} \sqrt{\frac{15}{2\pi}} \cdot \frac{(x + iy)^2}{r^2}$$

Recurrence Relation for Associated Laguerre Polynomials

$$\rho L_p^q(\rho) = (2p + q + 1)L_p^q(\rho) - \left[\frac{p + 1}{p + q + 1} \right] L_{p+1}^q(\rho) - (p + q)^2 L_{p-1}^q(\rho)$$

Orthonormality Condition for Electronic States of a Hydrogen Atom

$$\langle \varphi_{n'l'm'} | \varphi_{nlm} \rangle = \delta_{n'n} \delta_{l'l} \delta_{m'm}$$

4. Rydberg Atoms

The atoms excited to a high angular momentum state is called a Rydberg atom. Here, we explore the properties of the hydrogen atoms excited to the Rydberg state.

- a. Show that for a hydrogen atom in the state corresponding to maximum orbital angular momentum ($l = n - 1$), (10 points)

$$\langle n, n - 1 | r | n, n - 1 \rangle = a_0 n \left(n + \frac{1}{2} \right)$$
$$\langle n, n - 1 | r^2 | n, n - 1 \rangle = a_0^2 n^2 (n + 1) \left(n + \frac{1}{2} \right)$$

- b. Under the same conditions as in problem 4a, show that for large values of n and l , (10 points)

$$\begin{aligned}\sqrt{\langle \hat{r}^2 \rangle} &\rightarrow a_0 n^2 \\ \frac{\Delta r}{\langle r \rangle} &\rightarrow 0 \\ E_n &\rightarrow -\frac{1}{2} \frac{e^2}{n^2 a_0}\end{aligned}$$

where $(\Delta r)^2 = \langle \hat{r}^2 \rangle - \langle \hat{r} \rangle^2$.

5. Interacting spin systems (Cohen-Tannoudji, Liu and Laloe Exercise 2 in E_{XIII})

In this problem, we consider two spin $\frac{1}{2}$ particles coupled by an interaction Hamiltonian of the form $a(t)\vec{S}_1 \cdot \vec{S}_2$, where $a(t)$ approaches zero when $|t|$ approaches infinity, and takes on a non-negligible value on the order of a_0 only within the time interval $|t| \leq \tau$ near $t = 0$.

- a. At $t = -\infty$, the system is in the state $|\uparrow\downarrow\rangle$, where the first spin is in the $+$ eigenstate of \hat{S}_{1z} and the second spin is in the $-$ eigenstate of \hat{S}_{2z} . Without any approximation, calculate the state of the system at $t = +\infty$. Show that the transition probability $\mathcal{P}(\uparrow\downarrow \rightarrow \downarrow\uparrow)$ of finding the system in the $|\downarrow\uparrow\rangle$ state at $t = +\infty$ depends only on the integral $\int_{-\infty}^{\infty} a(t)dt$. (10 points)

- b. Calculate $\mathcal{P}(\uparrow\downarrow \rightarrow \downarrow\uparrow)$ by using first-order time-dependent perturbation theory. Discuss the validity conditions for such an approximation by comparing the results with that obtained in part a. (10 points)