

Density Matrices and Open Quantum Systems

ECE 590.01

Quantum Engineering with Atoms

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But first... Atom-light Interaction Recap!

Atom-light interactions:

$$\hat{H}_I = -\frac{q}{m} A_0 e^{-i\omega t} \exp(i\vec{k} \cdot \hat{\mathbf{r}}) \hat{\mathbf{e}} \cdot \hat{\mathbf{p}}$$

$$= -\frac{q}{m} A_0 e^{-i\omega t} (\boxed{1} + i\vec{k} \cdot \hat{\mathbf{r}} + \dots) \hat{\mathbf{e}} \cdot \hat{\mathbf{p}}$$

$[\hat{\mathbf{r}}, \hat{H}_0] = \frac{i\hbar}{m} \hat{\mathbf{p}}$

$$\hat{H}_I^{(\text{E1})} = -\frac{q}{m} \frac{E_0}{i\omega} e^{-i\omega t} \hat{\mathbf{e}} \cdot \hat{\mathbf{p}}$$

$$\hat{H}_I^{(\text{E1})} = \frac{q}{\hbar\omega} E_0 e^{-i\omega t} \hat{\mathbf{e}} \cdot (\hat{\mathbf{r}} \hat{H}_0 - \hat{H}_0 \hat{\mathbf{r}})$$

$\omega_{fi} \approx \omega$

$$\hat{H}_I = -q\vec{E}(0, t) \cdot \hat{\mathbf{r}}$$

Electric dipole (E1) interaction

First-order Interaction Term

$$\hat{H}_I = -\frac{q}{m}A_0e^{-i\omega t}(1 + \boxed{i\vec{k} \cdot \hat{\mathbf{r}}} + \dots)\hat{\mathbf{e}} \cdot \hat{\mathbf{p}}$$

$$H_I^{(1)} = -i\frac{q}{m}kA_0e^{-i\omega t}(\hat{\mathbf{k}} \cdot \hat{\mathbf{r}})(\hat{\mathbf{e}} \cdot \hat{\mathbf{p}})$$

$$(\vec{A} \cdot \vec{B})(\vec{C} \cdot \vec{D}) = (\vec{A} \times \vec{C}) \cdot (\vec{B} \times \vec{D}) + (\vec{B} \cdot \vec{C})(\vec{A} \cdot \vec{D})$$

$$H_I^{(1)} = -i\frac{q}{m}kA_0e^{-i\omega t} \left[\boxed{(\hat{\mathbf{k}} \times \hat{\mathbf{e}})(\hat{\mathbf{r}} \times \hat{\mathbf{p}})} + \boxed{(\hat{\mathbf{e}} \cdot \hat{\mathbf{r}})(\hat{\mathbf{k}} \cdot \hat{\mathbf{p}})} \right]$$

$$H_I^{\text{M1}} = -\frac{q}{m}\vec{B}(0,t) \cdot \hat{\mathbf{L}}$$

$$= -\frac{q}{m}\vec{B}(0,t) \cdot (\hat{\mathbf{L}} + g_s\hat{\mathbf{S}})$$

Magnetic dipole (M1) interaction

$$H_I^{(\text{E2})} = i\frac{qk}{\hbar\omega} \left(\vec{E}(0,t) \cdot \hat{\mathbf{r}} \right) \left[\hat{\mathbf{k}} \cdot (\hat{\mathbf{r}}\hat{H}_0 - \hat{H}_0\hat{\mathbf{r}}) \right]$$

$$= -iq(\vec{E}(0,t) \cdot \hat{\mathbf{r}})(\vec{k} \cdot \hat{\mathbf{r}})$$

Electric quadrupole (E2) interaction

Wave Functions vs Density Matrices

	Wave function	Density Matrix
Notation:	$ \psi(t)\rangle$	$\rho(t)$
Applies to:	Pure states	Pure or mixed states
Size:	$N \times 1$ vector	$N \times N$ matrix
Evolution:	Schrodinger equation	von Neuman equation
Superposition	$ \psi\rangle = \frac{1}{\sqrt{2}}(0\rangle + 1\rangle)$	$\rho_p = \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix}$ vs $\rho_m = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$

Density Matrix – General Form

$$\rho = \sum_i p_i |\psi_i\rangle \langle \psi_i|$$

Pure state: $p_i = \delta_{ik}$ for a single specific k

Normalization:

$$\begin{aligned}\text{Tr}(\rho) &= \sum_j \langle \psi_j | \rho | \psi_j \rangle \\ &= \sum_j \langle \psi_j | \sum_i p_i |\psi_i\rangle \langle \psi_i| \psi_j \rangle \\ &= \sum_i p_i \equiv 1\end{aligned}$$

Interpretation: p_i is probability to be in state $|\psi_i\rangle$

Density Matrix – Expectation Values

$$\begin{aligned}\langle \hat{O} \rangle &= \sum_i p_i \langle \psi_i | \hat{O} | \psi_i \rangle & \mathbb{1} &= \sum_j |\psi_j\rangle \langle \psi_j| \\ &= \sum_j \sum_i p_i \langle \psi_i | \psi_j \rangle \langle \psi_j | \hat{O} | \psi_i \rangle \\ &= \sum_i \sum_j p_j \langle \psi_i | \psi_j \rangle \langle \psi_j | \hat{O} | \psi_i \rangle \\ &= \sum_i \langle \psi_i | \left(\sum_j p_j |\psi_j\rangle \langle \psi_j| \hat{O} \right) | \psi_i \rangle \\ &= \sum_i \langle \psi_i | \rho \hat{O} | \psi_i \rangle = \text{Tr}(\rho \hat{O})\end{aligned}$$

$$\langle \hat{O} \rangle = \text{Tr}(\rho \hat{O}) = \text{Tr}(\hat{O} \rho) \quad (\text{Trace is cyclical})$$

Density Matrix – Pure vs Mixed States

$$\begin{aligned}\text{Tr}(\rho^2) &= \sum_k \langle \psi_k | \sum_j p_j |\psi_j\rangle \langle \psi_j | \sum_i p_i |\psi_i\rangle \langle \psi_i | \psi_k \rangle \\ &= \sum_k \sum_j \sum_i \delta_{kj} \delta_{ji} \delta_{ik} p_j p_i \\ &= \sum_i p_i^2 \quad \begin{array}{ll} = 1 & \Leftrightarrow p_i = \delta_{ik} \text{ (pure state)} \\ < 1 & \text{(mixed state)} \end{array}\end{aligned}$$

The trace of ρ^2 is necessary and sufficient to determine state purity.

In general: $\frac{1}{d} \leq \text{Tr}(\rho) \leq 1$ where d is the dimension of the Hilbert space

Density Matrix – Matrix Representation

$$\rho = \begin{pmatrix} \rho_{00} & \rho_{01} & \dots & \rho_{0N} \\ \rho_{10} & \rho_{11} & \dots & \rho_{1N} \\ \vdots & \vdots & \ddots & \vdots \\ \rho_{N0} & \rho_{N1} & \dots & \rho_{NN} \end{pmatrix}$$

$\rho_{ii} \in (\mathbb{R} \geq 0)$ are called the *populations* $\sum_{i=1}^N \rho_{ii} = 1$

$\rho_{ij} \in \mathbb{C}$ are called the *coherences* $\rho_{ij} = \rho_{ji}^*$

Density Matrix – Time Evolution

$$\rho = \sum_i p_i |\psi_i\rangle \langle \psi_i|$$

$$\frac{d\rho}{dt} = \sum_i p_i \left[\left(\frac{d}{dt} |\psi_i\rangle \right) \langle \psi_i| + |\psi_i\rangle \left(\frac{d}{dt} \langle \psi_i| \right) \right]$$

$$= \sum_i p_i \left[\frac{H}{i\hbar} |\psi_i\rangle \langle \psi_i| + |\psi_i\rangle \langle \psi_i| \frac{H}{-i\hbar} \right]$$

$$= -\frac{i}{\hbar} (H\rho - \rho H) = -\frac{i}{\hbar} [H, \rho] \quad (\text{von Neumann equation})$$

Integrate directly:

$$\rho(t) = e^{-iHt/\hbar} \rho(0) e^{iHt/\hbar}$$

Density Matrix – Time Evolution

$$\begin{aligned}\rho(t) &= \sum_i p_i |\psi_i(t)\rangle \langle \psi_i(t)| \\ &= \sum_i p_i e^{-iHt/\hbar} |\psi_i(0)\rangle \langle \psi_i(0)| e^{iHt/\hbar} \\ &= e^{-iHt/\hbar} \rho(0) e^{iHt/\hbar}\end{aligned}$$

Density Matrix – Partial Trace

$$\mathrm{Tr}_B(\rho_{AB}) = \rho_A \quad \mathcal{H}_A \otimes \mathcal{H}_B \longrightarrow \mathcal{H}_A$$

- Partial trace effectively averages over system B to yield the state of system A.
 - Typically used if we can't probe system B
- Must behave correctly: $\mathrm{Tr}(X\rho_A) = \mathrm{Tr}[(X \otimes \mathbb{1}_B)\rho_{AB}]$

$$\begin{aligned} \mathrm{Tr}[(X \otimes \mathbb{1}_B)\rho_{AB}] &= \sum_a \sum_b (\langle a| \otimes \langle b|)(X \otimes \mathbb{1}_B)\rho_{AB}(|a\rangle \otimes |b\rangle) \\ &= \sum_a \sum_b \sum_{b'} (\langle a| \otimes \langle b|)(X \otimes |b'\rangle \langle b'|)\rho_{AB}(|a\rangle \otimes |b\rangle) \\ &= \sum_a \sum_b \langle a| X \langle b|\rho_{AB}|b\rangle (|a\rangle) \\ &= \sum_a \langle a| X \rho_A |a\rangle = \mathrm{Tr}(X\rho_A) \end{aligned}$$

Open Quantum Systems – Problem Definition

Open quantum systems allow for coupling between the system of interest with an environment (bath)

- Bath often has infinite degrees of freedom
- Cannot perform measurements on bath
- No control over bath
- Goal is to describe system and eliminate bath dependence

$$\begin{array}{c}
 \mathbb{1} \otimes H_B \\
 \uparrow \\
 H_T = H + H_B + H_{SB} \\
 \downarrow \qquad \qquad \downarrow \\
 H \otimes \mathbb{1}_B \qquad \text{interaction}
 \end{array}$$

$$\begin{aligned}
 \dot{\rho}_T &= -\frac{i}{\hbar} \left[\overset{H_0}{\boxed{H + H_B}} + H_{SB}, \rho_{SB} \right] \\
 \tilde{\rho}_T &= e^{-iH_0 t/\hbar} \rho_T e^{iH_0 t/\hbar} \quad (\text{interaction picture})
 \end{aligned}$$

Open Quantum Systems – Interaction Picture

$$\begin{aligned}\dot{\tilde{\rho}}_T &= e^{iH_0t/\hbar} \left(i\frac{H_0}{\hbar}\rho_T + \dot{\rho}_T - i\rho_T\frac{H_0}{\hbar} \right) e^{-iH_0t/\hbar} \\&= -\frac{i}{\hbar} e^{iH_0t/\hbar} \left(-H_0\rho_T + [H_0 + H_{SB}, \rho_T] + \rho_T H_0 \right) e^{-iH_0t/\hbar} \\&= -\frac{i}{\hbar} e^{iH_0t/\hbar} (H_{SB}\rho_T - \rho_T H_{SB}) e^{-iH_0t/\hbar} \\&= -\frac{i}{\hbar} e^{iH_0t/\hbar} \left(H_{SB}e^{-iH_0t/\hbar} e^{iH_0t/\hbar} \rho_T - \rho_T e^{-iH_0t/\hbar} e^{iH_0t/\hbar} V \right) e^{-iH_0t/\hbar} \\&= -\frac{i}{\hbar} \left[\tilde{H}_{SB}, \tilde{\rho}_T \right]\end{aligned}$$

Open Quantum Systems – Time Evolution

$$\dot{\tilde{\rho}}_T = -\frac{i}{\hbar} \left[\tilde{H}_{SB}, \tilde{\rho}_T \right]$$

$$\tilde{\rho}_T(t) = \tilde{\rho}_T(0) - \frac{i}{\hbar} \int_0^t dt' \left[\tilde{H}_{SB}(t'), \tilde{\rho}_T(t') \right]$$

$$\dot{\tilde{\rho}}_T = -\frac{i}{\hbar} \left[\tilde{H}_{SB}, \tilde{\rho}_T(0) \right] - \frac{1}{\hbar^2} \int_0^t dt' \left[\tilde{H}_{SB}(t), \left[\tilde{H}_{SB}(t'), \tilde{\rho}_T(t') \right] \right]$$

$$\tilde{\rho}_S = \text{Tr}_B(\tilde{\rho}_T) \equiv \tilde{\rho} \quad + \text{system and bath start out uncorrelated}$$

$$\dot{\tilde{\rho}} = -\frac{1}{\hbar^2} \int_0^t \text{Tr}_B \left(\left[\tilde{H}_{SB}(t), \left[\tilde{H}_{SB}(t'), \tilde{\rho}_T(t') \right] \right] \right)$$

Open Quantum Systems – Approximations

Weak system-bath interaction
“Large” bath
(Born approximation)

$$\dot{\rho} = -\frac{1}{\hbar^2} \int_0^t \text{Tr}_B \left(\left[\tilde{H}_{SB}(t), \left[\tilde{H}_{SB}(t'), \tilde{\rho}_T(t') \right] \right] \right)$$

System kernel only depend
on current state
(Born-Markov approximation)

$$= -\frac{1}{\hbar^2} \int_0^t \text{Tr}_B \left(\left[\tilde{H}_{SB}(t), \left[\tilde{H}_{SB}(t'), \tilde{\rho}(t') \otimes \tilde{\rho}_B(0) \right] \right] \right)$$

$$= -\frac{1}{\hbar^2} \int_0^t \text{Tr}_B \left(\left[\tilde{H}_{SB}(t), \left[\tilde{H}_{SB}(t'), \tilde{\rho}(t) \otimes \tilde{\rho}_B(0) \right] \right] \right)$$

Open Quantum System – System-Bath Separation

$$\tilde{H}_{SB}(t) = \hbar \sum_i S_i(t) \otimes B_i(t) = \hbar \sum_i S_i^\dagger(t) \otimes B_i^\dagger(t) \quad (\text{Equivalent due to Hermiticity})$$

$$\dot{\tilde{\rho}} = -\frac{1}{\hbar^2} \int_0^t \text{Tr}_B \left(\left[\tilde{H}_{SB}(t), \left[\tilde{H}_{SB}(t'), \tilde{\rho}(t) \otimes \tilde{\rho}_B(0) \right] \right] \right)$$

$$\begin{aligned} \dot{\tilde{\rho}} = & - \sum_i \sum_j \int_0^t dt' \{ S_i^\dagger(t) S_j(t') \tilde{\rho}(t) \text{Tr}_B [B_i^\dagger(t) B_j(t') \tilde{\rho}_B(0)] \\ & - S_i(t) \tilde{\rho}(t) S_j^\dagger(t') \text{Tr}_B [B_i(t) \tilde{\rho}_B(0) B_j^\dagger(t')] \\ & - S_j(t') \tilde{\rho}(t) S_i^\dagger(t) \text{Tr}_B [B_j(t') \tilde{\rho}_B(0) B_i^\dagger(t)] \\ & + \tilde{\rho}(t) S_j^\dagger(t') S_i(t) \text{Tr}_B [\tilde{\rho}_B(0) B_j^\dagger(t') B_i(t)] \} \end{aligned}$$

Free to choose either substitution

Open Quantum Systems – Fast Bath Dynamics

$$\begin{aligned}\text{Tr} \left[B_j(t') \tilde{\rho}_B B_i^\dagger(t) \right] &= \text{Tr} \left[\tilde{\rho}_B B_i^\dagger(t) B_j(t') \right] = \text{Tr} \left[B_i^\dagger(t) B_j(t') \tilde{\rho}_B \right] \equiv \frac{h_{ij}}{2} \delta(t' - t) \\ \text{Tr} \left[B_i(t) \tilde{\rho}_B B_j^\dagger(t') \right] &= \text{Tr} \left[\tilde{\rho}_B B_j^\dagger(t') B_i(t) \right] = \text{Tr} \left[B_j^\dagger(t') B_i(t) \tilde{\rho}_B \right] \equiv \frac{h_{ji}}{2} \delta(t' - t)\end{aligned}$$

Correlation within
bath don't last

$$\begin{aligned}\dot{\tilde{\rho}} = & - \sum_i \sum_j \int_0^t dt' \{ S_i^\dagger(t) S_j(t') \tilde{\rho}(t) \text{Tr}_B \left[B_i^\dagger(t) B_j(t') \tilde{\rho}_B(0) \right] \\ & - S_i(t) \tilde{\rho}(t) S_j^\dagger(t') \text{Tr}_B \left[B_i(t) \tilde{\rho}_B(0) B_j^\dagger(t') \right] \\ & - S_j(t') \tilde{\rho}(t) S_i^\dagger(t) \text{Tr}_B \left[B_j(t') \tilde{\rho}_B(0) B_i^\dagger(t) \right] \\ & + \tilde{\rho}(t) S_j^\dagger(t') S_i(t) \text{Tr}_B \left[\tilde{\rho}_B(0) B_j^\dagger(t') B_i(t) \right] \} \end{aligned}$$

$$\dot{\tilde{\rho}} = - \sum_i \sum_j \left\{ \frac{h_{ij}}{2} \left[S_i^\dagger(t) S_j(t) \tilde{\rho}(t) - S_j(t) \tilde{\rho}(t) S_i^\dagger(t) \right] + \frac{h_{ji}}{2} \left[\tilde{\rho}(t) S_j^\dagger(t) S_i(t) - S_i(t) \tilde{\rho}(t) S_j^\dagger(t) \right] \right\}$$

Open Quantum Systems – First Lindblad Form

$$\begin{aligned}\dot{\tilde{\rho}} &= - \sum_i \sum_j \left\{ \frac{h_{ij}}{2} \left[S_i^\dagger(t) S_j(t) \tilde{\rho}(t) - S_j(t) \tilde{\rho}(t) S_i^\dagger(t) \right] + \frac{h_{ji}}{2} \left[\tilde{\rho}(t) S_j^\dagger(t) S_i(t) - S_i(t) \tilde{\rho}(t) S_j^\dagger(t) \right] \right\} \\ &= - \sum_i \sum_j \left\{ \frac{h_{ij}}{2} \left[S_i^\dagger S_j \tilde{\rho} - S_j \tilde{\rho} S_i^\dagger \right] + \frac{h_{ji}}{2} \left[\tilde{\rho} S_i^\dagger S_j - S_j \tilde{\rho} S_i^\dagger \right] \right\} \\ &= \sum_i \sum_j h_{ij} \left\{ S_j \tilde{\rho} S_i^\dagger - \frac{1}{2} \left(S_i^\dagger S_j \tilde{\rho} + \tilde{\rho} S_i^\dagger S_j \right) \right\}\end{aligned}$$

Open Quantum System – Standard Lindblad Form

$$\dot{\tilde{\rho}} = \sum_i \sum_j h_{ij} \left\{ S_j \tilde{\rho} S_i^\dagger - \frac{1}{2} \left(S_i^\dagger S_j \tilde{\rho} + \tilde{\rho} S_i^\dagger S_j \right) \right\}$$

Using an appropriate
unitary transform d :

$$\Gamma = d h d^\dagger \begin{cases} \Gamma_i = \sum_k \sum_l d_{ik} h_{kl} d_{li}^* \\ L_i = \sum_k d_{ik} S_k \end{cases}$$

$$\dot{\rho} = -\frac{i}{\hbar} [H, \rho] + \sum_i \Gamma_i \left[L_i \rho L_i^\dagger - \frac{1}{2} \left(L_i^\dagger L_i \rho + \rho L_i^\dagger L_i \right) \right]$$

$$\tilde{L}_i = \sqrt{\Gamma_i} L_i$$

$$\dot{\rho} = -\frac{i}{\hbar} [H, \rho] + \sum_i \left[\tilde{L}_i \rho \tilde{L}_i^\dagger - \frac{1}{2} \left(\tilde{L}_i^\dagger \tilde{L}_i \rho + \rho \tilde{L}_i^\dagger \tilde{L}_i \right) \right]$$